

# Dye-Encapsulated Zeolitic Imidazolate Framework (ZIF-71) for Fluorochromic Sensing of Pressure, Temperature, and Volatile Solvents

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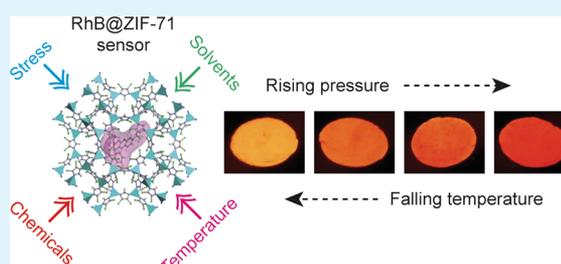
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**ABSTRACT:** Luminescent metal–organic frameworks (MOFs) offer a multifunctional platform for creating noninvasive sensors and tunable optoelectronics. However, fluorochromic materials that are photophysically resilient and show high sensitivity toward different physical and chemical stimuli are scarce. We report a facile host–guest nanoconfinement strategy to construct a fluorescent hybrid material with multiple sensing capabilities. We design and fabricate a new Guest@MOF material: comprising a zeolitic MOF (ZIF-71) as a nanoporous host for encapsulating rhodamine B (RhB dye) guest molecules, resulting in an RhB@ZIF-71 system with mechanochromic, thermochromic, and solvatochromic sensing response. The fluorochromic sensing properties stem from the nanoconfinement effect that ZIF-71 imposes on RhB monomers, yielding the H- or J-type aggregates with tunable photophysical and photochemical properties. For mechanochromism, the external pressure causes an emission red shift in a linear fashion, switching RhB guests from H-type to J-type aggregates through a shear deformation. For thermochromism, we demonstrate a linear scaling as a function of temperature due to the spatial restriction imposed on J-type aggregates incarcerated in ZIF-71 pores. Harnessing the solvatochromism of RhB@ZIF-71, we interrogated its photochemical response by employing three diverse groups of volatile organic compounds. The multimodal sensing response paved the way to smart applications like photonic pressure sensors, noninvasive thermometers, and ultrasensitive chemosensors.

**KEYWORDS:** hybrid materials, metal–organic frameworks, luminescence, guest–host interactions, photonic sensors



## 1. INTRODUCTION

Research on the luminescent sensing abilities of metal–organic framework (MOF) materials and their composite systems is rapidly expanding because of their wide range of technological applications in nanophotonics, bio-imaging, and smart sensors.<sup>1–4</sup> Compared to other luminescent sensing materials such as metal complexes<sup>5</sup> and dye-based fluorescent probes,<sup>6</sup> crystalline MOF exhibits long-range periodicity, nanosized pores, combined with diverse chemical and structural versatility for tuning its vast physicochemical properties.<sup>7</sup> Among the many reported luminescent MOFs,<sup>8</sup> the study of Guest@MOF systems<sup>9–14</sup> composed of an MOF structure (the “host”) to afford the encapsulation of luminescent molecules (the “guest”) has attracted considerable attention. This is because the Guest@MOF composite can be achieved by an easy-to-synthesize or “off-the-shelf” MOF structure and commonly available fluorophores (luminescent dyes) to yield bespoke luminescent properties.<sup>15–17</sup> Additionally, the MOF hosts will act as a shield, protecting the guest against possible sources of degradation, like in the case of the encapsulation of the unstable fluorescent perovskite nanocrystals, which enhanced their robustness when confined within MOFs, paving the way to their integration into functional photonic

devices.<sup>18–20</sup> Notably, the aforementioned approach could circumvent the need to rely on complex chemical designs to yield luminescent MOFs or to incorporate expensive rare-earth elements widely deployed in inorganic phosphors.<sup>21</sup>

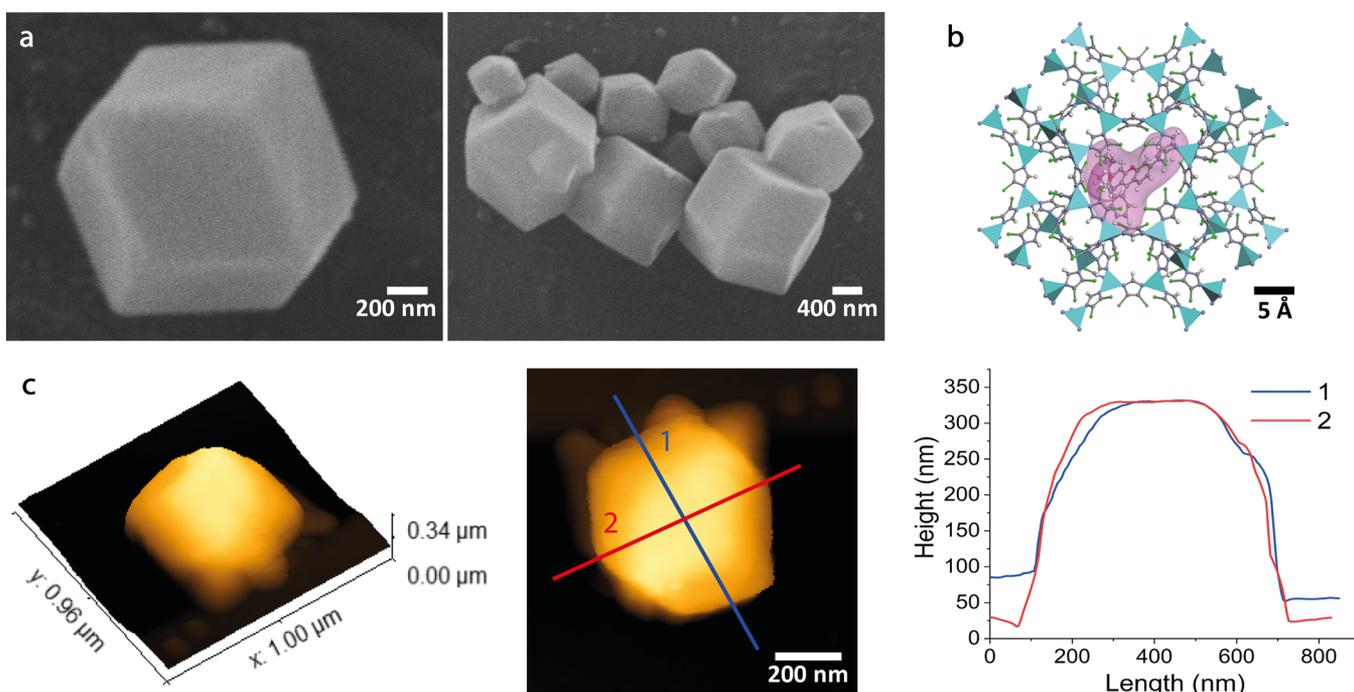
The existing Guest@MOF approach faces a number of challenges. For example, the vast majority of Guest@MOF sensing mechanisms are based on the luminescence quenching effect where the analytes “switch off” the emission of the guest.<sup>22–24</sup> However, this mechanism might be susceptible to moisture exposure or temperature fluctuation in the environment, giving inaccurate readouts.<sup>15,25</sup> There are other Guest@MOF materials exhibiting luminescent sensing through a “turn-on” sensing<sup>26</sup> or luminescence color (fluorochromic) switching mechanism, but they often require very specific chemical interactions to function,<sup>15</sup> thus limiting multimodal applica-

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**Figure 1.** (a) SEM micrographs of the submicron-sized crystals of RhB@ZIF-71 exhibiting a rhombic dodecahedron habit. (b) Illustration of the Guest@MOF framework structure of RhB@ZIF-71 where the pore of the ZIF-71 host is used to confine the luminescent RhB guest (represented by the molecule with a red surface). From the pore spatial constraint, it can be deduced that two RhB molecules could occupy a unit cell of ZIF-71. Color scheme: ZnN<sub>4</sub> tetrahedron in cyan, nitrogen in dark blue, carbon in gray, hydrogen in white, chlorine in green, and oxygen in red. (c) AFM height topography and the corresponding cross-sectional profiles of a representative RhB@ZIF-71 crystal.

tions. To address some of the foregoing challenges, we propose that fluorochromic sensing can be achieved by using MOF pores to induce different caging effects on luminescent dye monomers and their aggregates, thereby producing a distinctive fluorescent behavior when subject to different physical and chemical stimuli.

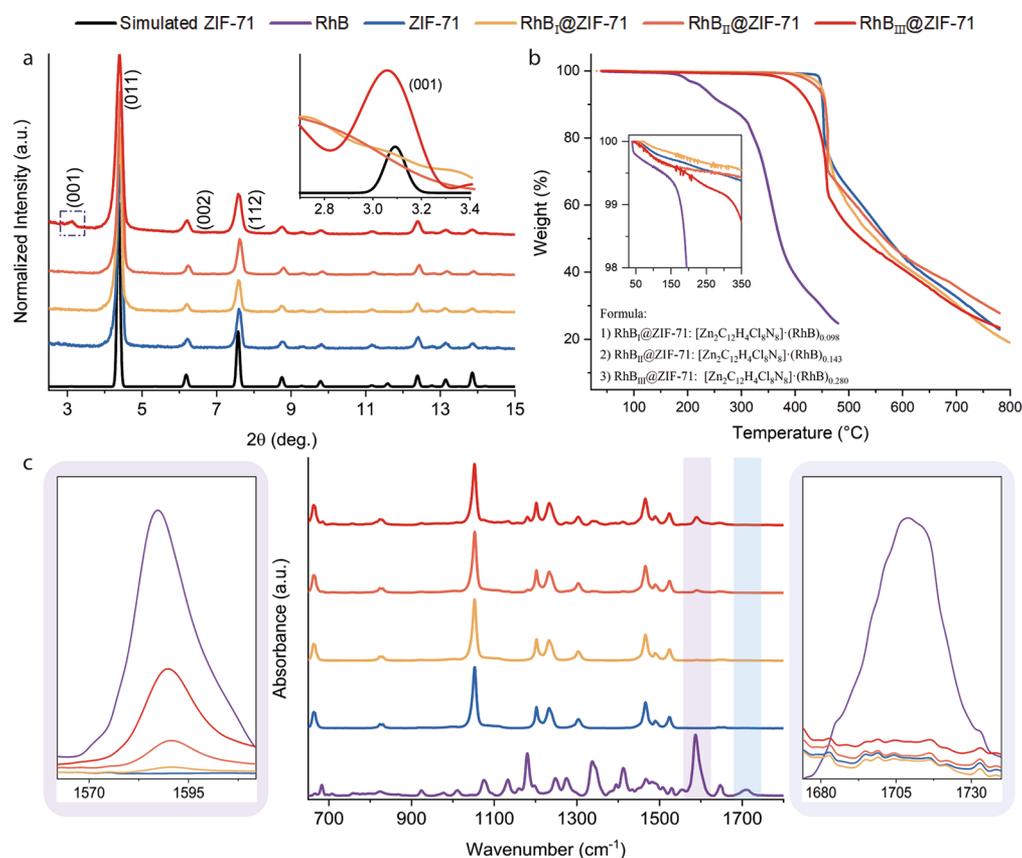
In this study, we demonstrate a facile method to synthesize a multi-stimuli responsive Guest@MOF luminescent material by embedding the well-known fluorescent dye rhodamine B (RhB) into zeolitic imidazolate framework-71 (ZIF-71). Although a number of RhB@MOF systems have been reported,<sup>27–34</sup> there are certain limitations, such as the need to use complex synthesis steps, no control of RhB aggregates, inadequate understanding of the underlying mechanisms, and single-mode sensing functionality (basically, only solvatochromism). In contrast, our RhB@ZIF-71 not only shows solvatochromic response with its described mechanism far surpassing all reported examples<sup>27–34</sup> but also exhibits intriguing mechanochromic and thermochromic sensing properties that are unknown to date. We employed detailed photophysical characterization techniques to unravel the underpinning guest–host mechanisms controlling the performance of this new composite system. Our results also demonstrate proof-of-concept RhB@ZIF-71 applications, opening new pathways for the realization of multimodal fluorochromic sensing platforms.

## 2. RESULT AND DISCUSSION

**2.1. Synthesis and Structure of the RhB@ZIF-71 System.** The RhB@ZIF-71 system was synthesized through rapid mixing of a solution of zinc acetate with a solution of 4,5-dichloroimidazole (dcIm) and RhB at room temperature; the full details are given in the [Experimental Section](#). Through this

straightforward one-step reaction, we observe that the mixed solution was immediately converted from a transparent pink color to turbid, indicating the rapid formation of RhB@ZIF-71. As shown in [Figure 1](#), the crystal size of RhB@ZIF-71 was found to be ~800 nm by scanning electron microscopy (SEM) and atomic force microscopy (AFM). We also characterized the morphology of RhB@ZIF-71 at different reaction times ([Figure S1](#), Supporting Information) and found that, after ~10 min, a crystal size of ~800 nm has formed. The rapid formation of such submicron-sized crystals is consistent with the previous reports on pure ZIF-71.<sup>35</sup>

Powder X-ray diffraction (PXRD) was carried out to verify the crystal structure. As shown in [Figure 2a](#) and [Figure S2](#), the PXRD patterns of RhB@ZIF-71 with different RhB concentrations and pure ZIF-71 were consistent with the simulated XRD pattern of ZIF-71, indicating that the introduced guest (RhB) does not affect the structure of ZIF-71 significantly. However, when a very high concentration of RhB (0.5 mmol) was used, we observed the appearance of the (001) peak at  $2\theta = 3.1^\circ$  ([Figure 2a](#), inset), while this peak was not obvious in the samples whose RhB concentrations are relatively low (0.01 and 0.05 mmol). Although the peak at  $2\theta = 3.1^\circ$  exists in the simulated pattern, its relative intensity is typically too low to be observed by XRD. Since a smaller  $2\theta$  value represents a larger crystal plane separation, the appearance of the (001) peak suggests that many ZIF-71 pores may contain more than one RhB molecule, thereby affecting the preferred orientation during crystal growth. Because the size of one RhB molecule is around  $16.19 \times 12.83 \times 6.97$  Å (including the van der Waals surface) and the minimum/maximum distance inside the ZIF-71 pore is 16.58/22.59 Å, it is conceivable that the crystal growth preference will be influenced when more than one RhB molecule is stacked together. This is further supported by the



**Figure 2.** (a) Comparison of the XRD patterns of the simulated/synthesized ZIF-71 and RhB@ZIF-71 featuring three RhB concentrations. Inset: the (001) peak of simulated ZIF-71 is enlarged to compare with RhB<sub>III</sub>@ZIF-71, while this peak is absent for both RhB<sub>I</sub>@ZIF-71 and RhB<sub>II</sub>@ZIF-71. The simulated pattern of ZIF-71 was generated from the crystallographic information file (CIF) obtained from the Cambridge Structural Database (CCDC code: GITVIP). (b) TGA and (c) ATR-FTIR results for ZIF-71, RhB, and RhB@ZIF-71. Subscripts I, II, and III designate the concentration of RhB used in the synthesis to be 0.01, 0.05, and 0.5 mmol, respectively. The method to determine the chemical formulae derived from TGA data is presented in Section S1.

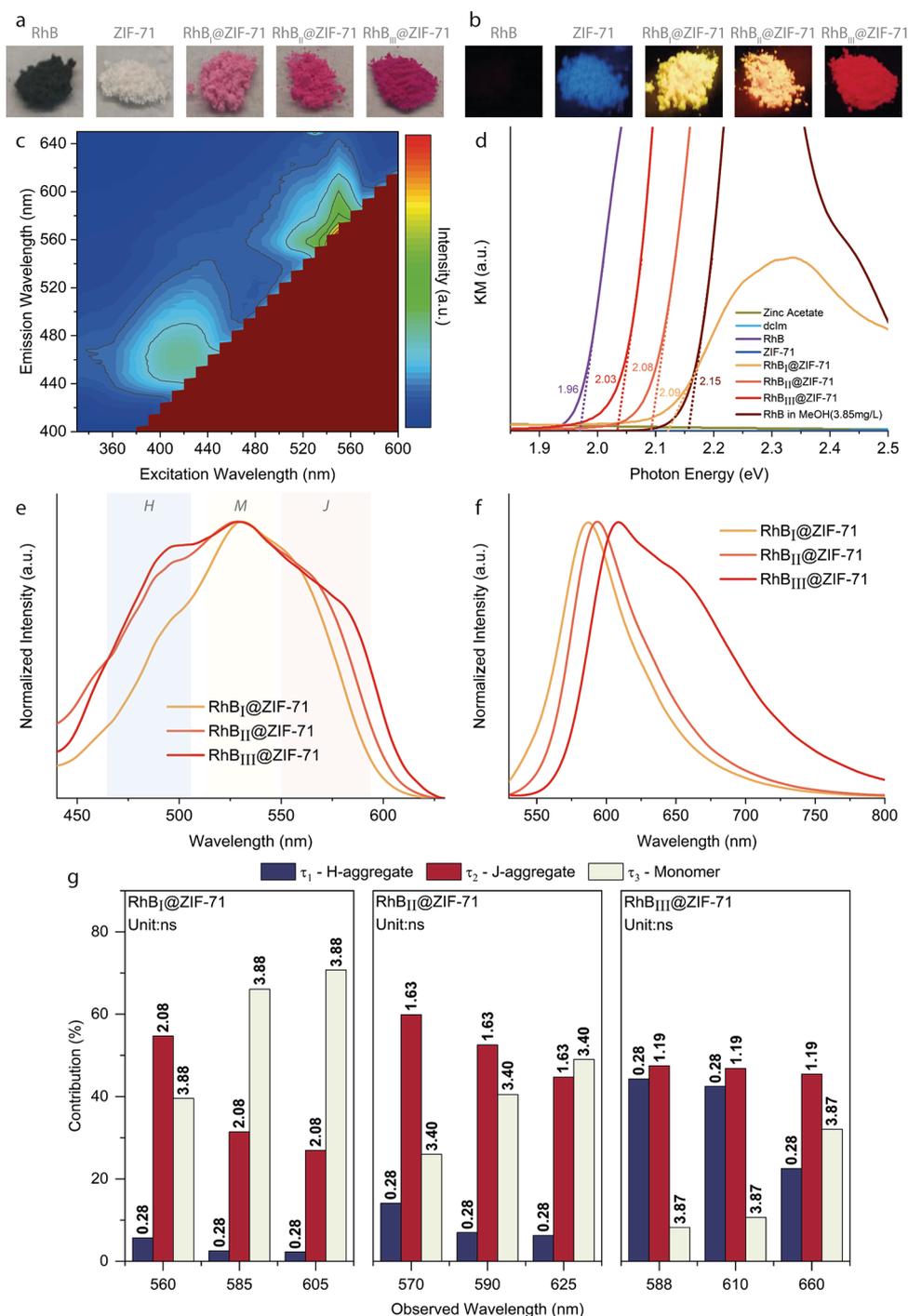
results of thermogravimetric analysis (TGA) (Figure 2b) from which the determined chemical formula shows there was on average more than one RhB per pore.

Although we did not see the trace of RhB in the PXRD patterns (Figure S2), it was successfully detected by Fourier transform infrared spectroscopy with attenuated total reflection (ATR-FTIR) (Figure 2c), TGA (Figure 2b), and Raman vibrational spectroscopy (Figure S3). Figure 2b,c demonstrates that the concentration of RhB inside the RhB@ZIF-71 system gradually increases as the amount of RhB used in the synthesis was increased, as evidenced in the systematic rise of the  $\sim 1590$  cm<sup>-1</sup> mode shown in the left inset of Figure 2c. Intriguingly, from the right inset of Figure 2c, we established that the peak at  $\sim 1710$  cm<sup>-1</sup> attributing to the C=O vibrational mode of RhB has completely disappeared after confinement. Likewise, we detected similar phenomena in the Raman spectra (Figure S3). Vibrational spectroscopic data revealed that the C=O of RhB could interact with the zinc atoms of ZIF-71 or with its dcIm linkers. The TGA results (Figure 2b) not only enabled us to derive the chemical formula of RhB@ZIF-71 but also revealed the enhanced thermal stability of RhB when confined in the system; this finding supports the notion that RhB guests are residing in the pores of ZIF-71.

**2.2. Luminescent Properties of the RhB@ZIF-71 System and Its Constituents.** We begin by investigating the emission properties of the pristine crystals of ZIF-71 and its dcIm linker under room temperature, as depicted in Figure

3a,b. It can be seen in Figure S4 that the dcIm linker displays an intense and broadband emission in the solid state with an emission maximum at  $\sim 468$  nm (under a 360 nm UV excitation), which can be attributed to the  $\pi^*-\pi$  transition.<sup>36</sup> From the emission map (Figure 3c) and the emission spectra (Figure S4) of ZIF-71, we established that ZIF-71 exhibits two emission peaks at around 456 and 559 nm. Because neither dcIm nor Zn(II) has an emission at 559 nm, the 559 nm emission of ZIF-71 could be from the ligand–metal charge transfer (LMCT), and the 456 nm emission comes from the linker itself in ZIF-71. Previous literature<sup>37</sup> mentioned that the LMCT process is usually expressed in MOF containing Zn(II), especially when the linker contains benzene derivatives, and MOFs often emit a green color fluorescence (500–565 nm) when LMCT occurs. Clearly, the structure and performance of ZIF-71 almost completely conform to this commonality, which supports our reasoning.

To confirm the aforementioned phenomena, we measured the emission lifetimes of dcIm (Table 1) and ZIF-71 (Table 2) employing the time-correlated single-photon-counting (TCSPC) technique. The most noticeable variation is when the observed wavelength changed from 450 to 558 nm (in Table 2); the  $\tau_3$  of ZIF-71 dramatically increased, which we attribute to the effect of LMCT. However, it is not certain that  $\tau_3$  is the lifetime of LMCT because the  $\tau_3$  of dcIm itself is 4.44 ns (in Table 1), and after forming the ZIF-71 framework structure, theoretically, it will increase due to the caging



**Figure 3.** (a) RhB, ZIF-71, and RhB@ZIF-71 with different RhB concentrations seen in the visible light and (b) their luminescence under the 365 nm UV excitation. (c) Emission map of the ZIF-71 powder. (d) Kubelka–Munk (KM) function for estimating the band gaps based on the photon energy intercepts. (e) Normalized excitation spectra (measured under em@650 nm) and (f) the normalized emission spectra (measured under ex@515 nm). (g) Lifetime data of RhB@ZIF-71 obtained using three different RhB concentrations.

effect.<sup>38</sup> The increase in  $\tau_3$  of dclm itself may be very close to the lifetime of the LMCT, which may cause the two to become indistinguishable. Thus, we consider that the  $\tau_3$  of dclm itself and the lifetime of LMCT together constitute the  $\tau_3$  of ZIF-71.

Subsequently, we characterized the band gap (Figure 3d), solid-state excitation and emission spectra (Figure 3e,f), absorption (Figure S5) of RhB@ZIF-71, and the emission spectra (Figure S6) of pure RhB solution with different concentrations at room temperature. However, the emission of

all the RhB@ZIF-71 powders is dominated by the guest itself (rather than simultaneously manifested by emissions of the guest and the LMCT of ZIF-71). The other possible interactions involving C=O with the zinc atoms or the linkers or the possible interaction between nitrogen atoms/the xanthene ring of RhB and the open metal sites/linker<sup>39</sup> may interrupt the LMCT process causing the single emission peak of RhB@ZIF-71. Of course, we do not rule out the possibility

**Table 1. Values of Time Constants ( $\tau_i$ ), Normalized Pre-Exponential Factors ( $a_i$ ), and Fractional Contributions ( $c_i = \tau_i \cdot a_i$ ) of the Emission Decay of the Pristine dcIm Linker in a Solid State and in Methanol Solutions upon Excitation at 362.5 nm ( $R_t = \sum a_i e^{(-t/\tau_i)}$ ,  $R_t$  is the Quantity/Counts at Time  $t$ )**

dcIm	$\lambda_{\text{obs}}$ [nm]	$\tau_1$ [ns]	$a_1$	$c_1$ [%]	$\tau_2$ [ns]	$a_2$	$c_2$ [%]	$\tau_3$ [ns]	$a_3$	$c_3$ [%]	$\chi^2$
solid state	470	0.36	0.041	11.89	1.76	0.038	53.95	4.44	0.010	34.16	1.052
0.0365 M	442	0.36	0.057	17.57	1.76	0.033	49.49	4.44	0.009	32.94	1.115
0.5 M	423	0.36	0.061	20.28	1.76	0.034	55.09	4.44	0.006	24.63	1.258
	443	0.36	0.060	18.80	1.76	0.037	55.64	4.44	0.007	25.55	1.082
	463	0.36	0.058	18.26	1.76	0.036	54.73	4.44	0.007	27.02	1.168

**Table 2. Values of Time Constants ( $\tau_i$ ), Normalized Pre-Exponential Factors ( $a_i$ ), and Fractional Contributions ( $c_i = \tau_i \cdot a_i$ ) of the Emission Decay of the Pristine ZIF-71 Powder upon Excitation at 362.5 nm**

ZIF-71	$\lambda_{\text{obs}}$ [nm]	$\tau_1$ [ns]	$a_1$	$c_1$ [%]	$\tau_2$ [ns]	$a_2$	$c_2$ [%]	$\tau_3$ [ns]	$a_3$	$c_3$ [%]	$\chi^2$
solid state	450	0.58	0.049	26.43	2.11	0.029	57.24	5.61	0.003	16.32	1.150
	558	0.58	0.041	14.74	2.11	0.026	34.50	5.61	0.015	50.76	1.215

that the emission of the ZIF-71 host could not be observed due to its relatively low quantum yield (Table S1).

In terms of their emission spectra (Figure 3f), it is shown that the bathochromic (red) shift was detected when the concentration of RhB increased. We propose that one of the reasons for this red shift is related to the formation of more RhB aggregates. Because of the relative dimensions of RhB molecules and ZIF-71 pores, the XRD results (Figure 2a) and the chemical composition derived from TGA results (Figure 2b) suggest that more than one RhB molecule may occupy the pore of ZIF-71. It follows that more RhB introduced during synthesis will lead to more aggregates. In principle, both of H-type (head-to-head) and J-type (head-to-tail) aggregates of RhB are able to form during the synthesis.<sup>40,41</sup> Since the emission of H-type aggregates is theoretically forbidden and the emission of J-type aggregates is allowed but with a longer wavelength,<sup>42</sup> we could observe the red shift. This idea is confirmed by the excitation spectra in Figure 3e. It can be seen that there are three excitation areas: humps at around 495 nm, peaks at 531 nm, and shoulders located between 552 and 590 nm. Using Kasha's exciton model,<sup>42</sup> we assigned the humps to H-type aggregates, peaks to RhB monomers, and the shoulders to J-type aggregates. Compared with the peaks (~531 nm), the intensity of the humps (~495 nm) increased so it follows that the amount of aggregates increased when a higher concentration of RhB was used during synthesis. The reason why the shoulders (552–590 nm) did not increase is explained below.

Comparing the variation of H-type and J-type aggregates, we found that the excitation peaks of the J-type aggregates showed a red shift as the concentration of RhB increases, while the H-type did not show a blue shift. This difference arises because the J-type aggregates have a longer spatial dimension than the H-type aggregates,<sup>41</sup> which might allow the J-type aggregates to interact with another J-type aggregate in adjacent pores. Our group has reported that perylene@ZIF-8 also showed a similar phenomenon in which perylene and 2-methylimidazole can form an energy transfer pathway through the adjacent pores.<sup>43</sup> Likewise, other researchers have demonstrated the preparation of long-range crystalline MOFs by mechanochemistry under solid conditions,<sup>44,45</sup> which means that chemical reactions can occur between solid crystals. On this basis, weak interactions involving J-aggregate interaction across the pores are plausible. Moreover, because the maximum size of the ZIF-71 window aperture is 5.08 Å, which is spatially larger than some parts of the RhB molecules (e.g., the distance of C–C on the xanthen

ring is 4.79 Å), we suggest that part of the J-type aggregates may protrude out of the ZIF-71 window, which will strengthen the interaction by bridging the pores. In contrast, the packing of H-type aggregates is tighter, and the occupied space is relatively small,<sup>41</sup> allowing them to be better confined inside the pores and less likely to interact with guests in adjacent pores. Hence, when the concentration of RhB used in the synthesis was increased, it will cause more J-type aggregate-based interactions (bridging the pores) in the RhB@ZIF-71 system, which leads to a smaller band gap (Figure 3d). This is another reason for the observed red shift in the emission and the reason behind the declining intensity of the emission shoulders (552–590 nm in Figure 3e).

To further study the formation of aggregates and analyze their luminescent properties, the emission lifetimes of pure RhB and RhB@ZIF-71 were measured by the TCSPC technique. For pure RhB in MeOH solution (0.0001 M), the lifetime is 2.93 ns, and for all the RhB@ZIF-71 powders, we obtained three decay times, as summarized in Figure 3g (see also Table S2 and Figure S7). We propose that  $\tau_1$  can be assigned to the H-type aggregates,  $\tau_2$  corresponds to the J-type aggregates, and  $\tau_3$  is due to the RhB monomers. First, it can be seen that  $\tau_3$  is greater than 2.93 ns, which also is evidence that RhB is residing inside the pore because the vibration of RhB monomers confined in the pore becomes restricted (i.e., caging effect), greatly reducing their nonradiative decay and thus increasing the lifetime. Second, it can also be seen that, as the RhB concentration increases, the values of  $a_1$  and  $c_1$  rose, while  $a_3$  and  $c_3$  fell (Table S2), which indicates an increase in the content of H-type aggregates and a relative decrease in the monomer content. Third,  $\tau_2$  was found to decrease as the RhB concentration increases; this supports the notion that the J-type aggregates could interact with each other across the adjacent pores. This kind of interaction can also be proven by comparing  $a_2$  and  $c_2$  (Table S2). As the observed wavelength ( $\lambda_{\text{obs}}$ ) increases, the magnitude of the change in  $a_2$  and  $c_2$  decreases, which results in more interactions and hence broadening of the emission component of the J-type aggregation. However, comparing the changes in  $a_1$ ,  $c_1$ ,  $a_2$ , and  $c_2$ , we can see that the increase of J-type aggregates is not as high as the H-type aggregates. Given the space limitation of the ZIF-71 pore, the larger J-type aggregates are harder to form than the H-type. Herein, the lifetime data have substantiated the previous inferences obtained from the emission and excitation spectrum.

The quantum yields of RhB<sub>I</sub>@ZIF-71, RhB<sub>II</sub>@ZIF-71, and RhB<sub>III</sub>@ZIF-71 were characterized, and the results are summarized in Table 3.

**Table 3. Quantum Yield of RhB@ZIF-71 with Different RhB Concentrations**

sample	QY <sup>a</sup> [%]		QY <sup>b</sup> [%]	
	ex@485 nm	ex@525 nm	ex@485 nm	ex@525 nm
RhB <sub>I</sub> @ZIF-71	23.99	28.25	35.43	39.53
RhB <sub>II</sub> @ZIF-71	13.74	14.15	17.51	18.30
RhB <sub>III</sub> @ZIF-71	1.68	1.85	2.32	3.31

<sup>a</sup>Samples were directly measured. <sup>b</sup>Samples (10 wt %) were first mixed with BaSO<sub>4</sub> (90 wt %) and then measured.

**2.3. Mechanochromic Sensing Response.** To study the mechanochromism of the RhB@ZIF-71 system, the RhB<sub>II</sub>@ZIF-71 material was chosen and compressed into pellets under different pressures. Figure 4a,b depicts the color of the pellets viewed under daylight and their emissions when subject to a 365 nm UV excitation, respectively. Here, we focus on the RhB<sub>II</sub>@ZIF-71 pellets because the RhB<sub>III</sub>@ZIF-71 has a relatively low quantum yield (Table 3), while the emission of the RhB<sub>I</sub>@ZIF-71 pellets (Figure S8) is not as linear as RhB<sub>II</sub>@ZIF-71. Figure 4d and Figure S9b reveal that the emission spectra of these pellets red-shifted as the pressure was systematically increased up to ~350 MPa, demonstrating a very linear relationship that is highly desirable for stress sensing applications (Figure 4f).

Combined with the excitation spectra (Figure 4c and Figure S9a), the lifetime data (Figure 4g and Figure S10 and Table S3), and the PXRD patterns (Figure 4e and Figure S11) of the pellets, we investigate the reason for the observed red shift in the emission. On the one hand, it can be seen that the excitation peaks of J-type aggregates (575–590 nm) showed a red shift with increasing pressure, and in the PXRD pattern (Figure 4e and Figure S11), it can be seen that the pressure has caused some amorphization of the ZIF-71 structure. These results suggest that the pelleting pressure leads to the mechanical deformation of the ZIF-71 structure where framework distortion will cause the tighter packing of aggregates and make the adjacent pores to come closer, causing stronger interactions inside the pores and promoting stronger interactions between the J-type aggregates across adjacent pores. Together, these factors result in a red shift. In principle, the stronger the interaction, the shorter the luminescent lifetime becomes,<sup>46</sup> but the  $\tau_2$  in Figure 4g did not decrease. This is because RhB aggregates are present in the pores of ZIF-71, and hence, the pore shrinkage from mechanical stress introduces a stronger caging effect, suppressing the nonradiative decay and preventing the decrease of lifetime. Based on this hypothesis, it is easy to understand the increase of the monomer's lifetime inside the pores ( $\tau_3$ ), which is dependent only upon the caging effect.

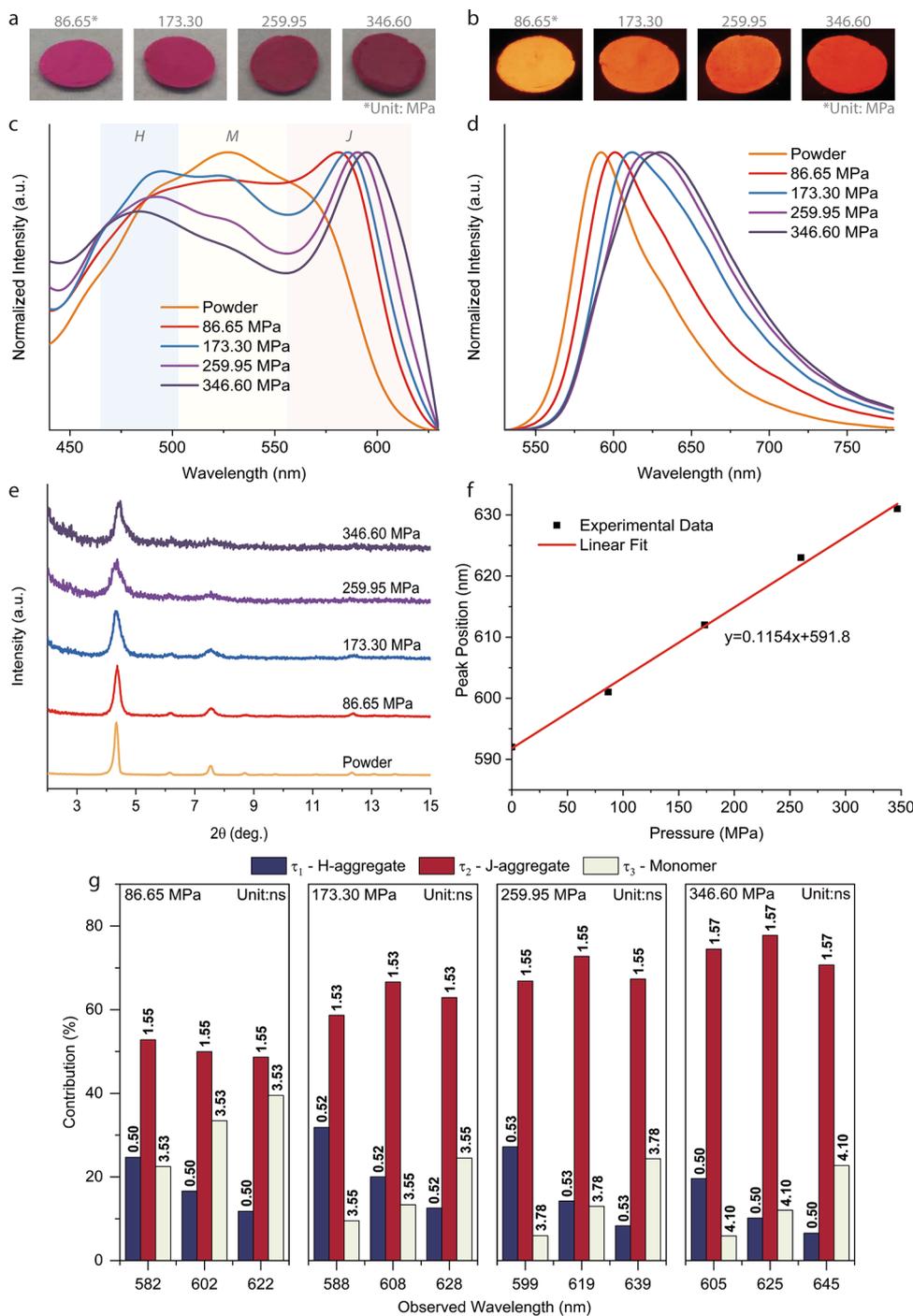
On the other hand, we consider that the red shift of the emission is also related to the increase in the relative content of the J-type aggregates. Table S4 shows that the RhB@ZIF-71 pellets had a smaller FWHM (full width at half-maximum) than the ZIF-71 pellets under pressure, which means that the RhB@ZIF-71 pellets possess a higher crystallinity than ZIF-71 (Figure S12) and reveals that the encapsulated guests can mechanically enhance the structural stability of the overall

framework under stress. Therefore, we propose the reason why  $a_3$  and  $c_3$  in Table S3 started to drop significantly at a relatively low pressure,  $a_1$  and  $c_1$  rose first and then fell, and  $a_2$  and  $c_2$  continuously increased is that the pores containing monomers are initially destroyed to form new H-type and J-type aggregates at a relatively low pressure. Subsequently, as the pressure keeps rising, the ZIF-71 crystals continue to deform under shear deformation,<sup>47</sup> causing the aggregates to transform from H-type to J-type (Scheme 1). Moreover, in the excitation spectra of the RhB<sub>II</sub>@ZIF-71 pellets (Figure 4c), the intensity of the H-type aggregates first rose and then fell, but the peak position of H-type aggregates (450–506 nm) is always blue-shifted with increasing pressure, which also supports this hypothesis.

**2.4. Thermochromic Sensing Response.** RhB@ZIF-71 exhibits thermochromic behavior as a function of temperature; here, we use RhB<sub>II</sub>@ZIF-71 as an example to explain the underlying mechanism. Figure 5a,b shows the excitation and emission spectra from room temperature to 200 °C, respectively. As shown in Figure 5 and Figure S13, the luminescent intensity of RhB@ZIF-71 decreases, accompanied by a red shift with increasing temperature. The decrease in intensity is very similar to the performance of pure RhB itself in solution at different temperatures,<sup>48,49</sup> which is due to the increase in a nonradiative decay rate. In the previous research,<sup>48</sup> the temperature range of this kind of intensity decrease of RhB was generally from 5 to 80 °C. Above this temperature range, the luminescent intensity of RhB was too low to be detected accurately. Remarkably, we demonstrate that the nonradiative decay of RhB was greatly reduced due to the caging effect of ZIF-71; not only overcoming the restriction that thermochromism of RhB can only be achieved in solutions, but our solid-state system also significantly extends the operational temperature range by at least a factor of 2 (Figure 5d). Figure S13d shows that, when the temperature rises from 200 to 250 °C, RhB@ZIF-71 experienced a relatively large red shift, which can be attributed to the thermal decomposition of RhB itself beyond 200 °C (consistent with TGA results in Figure 2b).

By analyzing the temperature range from room temperature to 200 °C, we established that there is a linear scaling relationship between the red shift of the emission peak and the temperature increment (Figure 5d); this effect is highly attractive for photonics-based thermometry applications. Note that the red shift of RhB in the solid state observed here is as yet unreported in the literature, unlike for pure RhB solutions in which the increase in temperature has produced no red shift.<sup>48</sup> The normalized excitation spectra (Figure 5a) reveal that, in this linearly changing region (room temperature to 200 °C), the relative intensity of H-type aggregates and monomers did not change much with the increase of temperature, while the J-type aggregates showed a relatively large intensity enhancement at high temperatures (i.e., 150 and 200 °C). Since the crystal structure of RhB@ZIF-71 can withstand a temperature of up to 250 °C (evidenced from the XRD patterns in Figure 5c and Figure S13e), we reason that the different spatial sizes between the H- and J-type aggregates and its monomers determine the trend of their intensity and wavelength change.

As discussed above, H-type aggregates and monomers possess smaller size than J-type aggregates, so the J-type aggregates will experience a stronger caging effect, therefore becoming less sensitive to temperature variation. In other



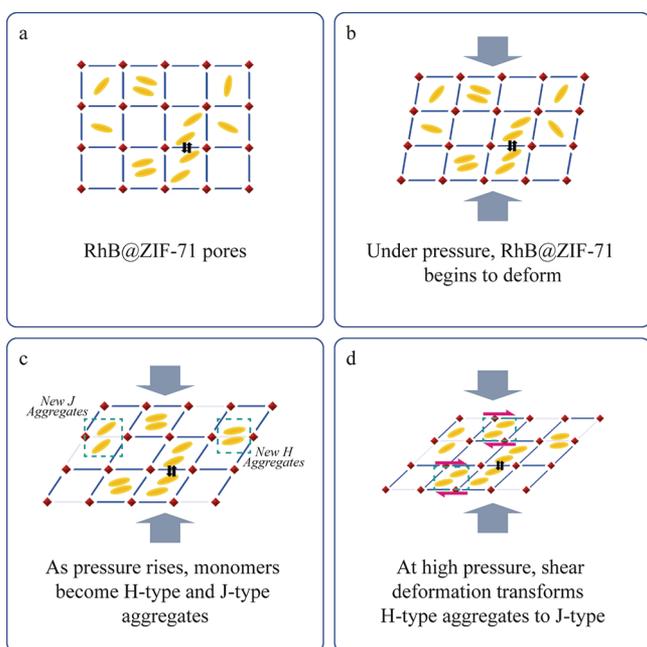
**Figure 4.** (a) RhB@ZIF-71 pellets prepared using different pelleting pressures, their colors viewed in visible light, and (b) their luminescence under the 365 nm UV excitation. (c) Normalized excitation spectra (measured under em@650 nm) and (d) normalized emission spectra (measured under ex@515 nm) of the RhB@ZIF-71 pellets. (e) XRD of the RhB@ZIF-71 pellets. (f) Linear relationship between the emission peak wavelength and the applied pressure for RhB@ZIF-71. (g) Lifetime data of RhB@ZIF-71 pellets showing the contributions from the monomer, H-aggregates, and J-aggregates.

words, when temperature rises, the excitation intensity of J-type aggregates decreases less than that of H-aggregates and monomers. It is this effect that causes the red shift observed in the emission peak. Additionally, at high temperatures, it can be seen that the excitation peaks of the H-type and J-type aggregates showed a very small degree of red and blue shifts, respectively, but the peaks associated with the monomers remain unchanged. This indicates that both H-type and J-type aggregates slightly expand at high temperatures, thereby

weakening the interactions within themselves. The absence of any variation to the excitation peaks of monomers reveals that the interactions between RhB and ZIF-71 are weak or negligible. Likewise, in the mechanochromism studies, we note that the peak wavelength of the monomer also did not change when subject to mechanical stress (Figure 4c), which also indicates that the interactions are weak or negligible.

**2.5. Solvatochromic Sensing Response.** Another promising property of RhB@ZIF-71 is its solvatochromic

### Scheme 1. Proposed Deformation of the RhB@ZIF-71 Subject to a Mechanical Pressure<sup>a</sup>



<sup>a</sup>Red nodes and blue rods represent the  $ZnN_4$  coordination environment and the dclm linker, respectively. Yellow ellipses represent the RhB guest molecules within the pore of the ZIF-71 host. (a) Possible interactions between J-aggregates are denoted by a pair of black arrows. (b, c) Illustration of how the structural distortion under pressure causes the formation of new H- and J-type aggregates. The gray rods represent the broken linkers under pressure. (d) Due to shear deformation (denoted by antiparallel pink arrows), the majority of the H-type aggregates transform into the J-type.

response, as shown in Figure 6 and Figure S14. To understand the subtle changes of emission due to solvatochromism, we chose the RhB<sub>1</sub>@ZIF-71 system, which has the lowest RhB concentration. We observed that the peak intensity and peak position of RhB<sub>1</sub>@ZIF-71 were distinctively different when exposed to different solvents. To explain this phenomenon, we propose that this solvatochromic response is linked to the nature of RhB itself. Because RhB can exist in three different forms in the solution state (i.e., lactone, zwitterion, and cation),<sup>48</sup> the lactone has no color under visible light or UV, while the zwitterion and cation have luminescence but their emission wavelength and intensity are different. Generally, polar protic solvents can stabilize the zwitterion<sup>48</sup> and therefore will give an intense luminescence. Meanwhile, in highly polar aprotic solvents, such as dimethylformamide (DMF), RhB has no luminescence due to the complete conversion to lactone, and in less polar aprotic solvents, such as acetonitrile (ACN), it can show luminescence.<sup>49</sup> Thereby, RhB is solvatochromic,<sup>48,50</sup> and we propose that this is also the basis for the solvatochromism observed in the RhB@ZIF-71 system.

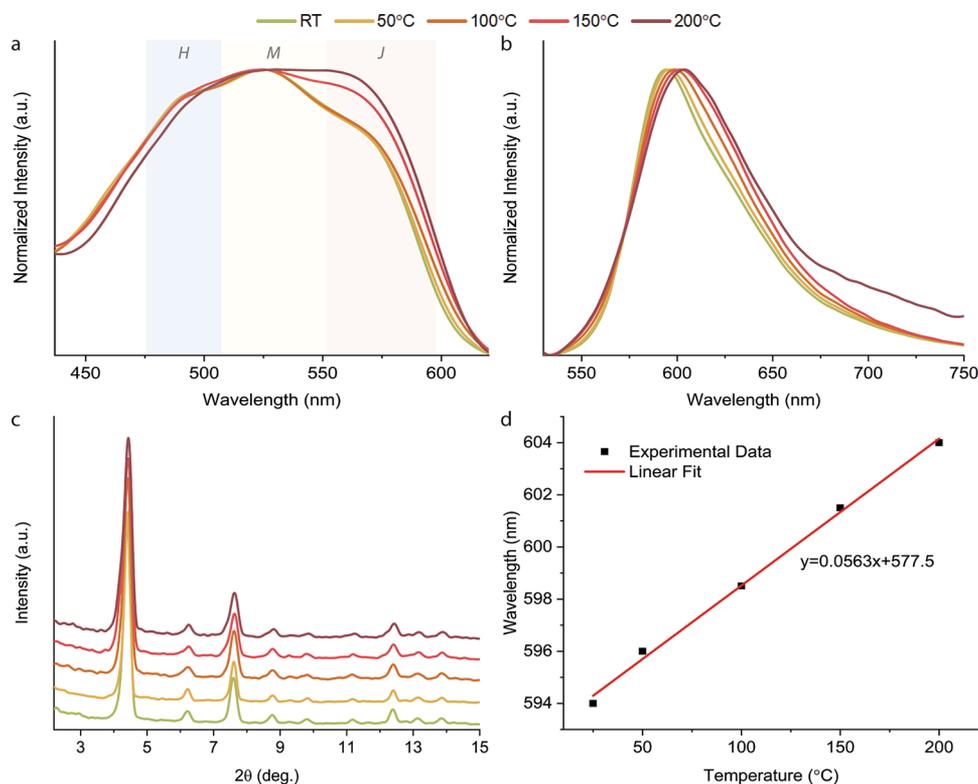
However, after confinement within ZIF-71, we found that the luminescence of RhB@ZIF-71 is greatly enhanced compared with the pure RhB (Figure S14). Notably, pure RhB has poor solubility or it is simply insoluble in many nonpolar solvents. For example, in this study, we tested toluene, hexane, and cyclohexane and confirmed that there was no luminescence in these hydrocarbons due to the

mentioned limitations. Conversely, we discovered that the RhB@ZIF-71 crystals exhibit good dispersion and luminescence (Figure S14 and Figure 6) in hexane and also in cyclic hydrocarbons (e.g., toluene and cyclohexane), thereby demonstrating a sensing property previously not achievable by pure RhB alone. To date, most of the RhB@MOF studies<sup>27–34</sup> have focused on the field of solvatochromism, and many of these systems possess a similar sensing behavior. However, almost all the authors attributed this kind of sensing to different solvents that affect the energy transfer between the MOF they used and RhB and lack of in-depth understanding of the emission wavelength change. Moreover, their theory is also unable to satisfactorily explain the luminescence of RhB@MOF systems in some highly polar aprotic solvents because no matter how the energy transfer occurs, when RhB turns into the lactone in strong polar aprotic solvents, in principle, the luminescence of the systems will be largely deteriorated.

Here, our systematic analysis of RhB@ZIF-71 based on the new confinement strategy provides new insight into the complex solvatochromism of RhB@MOF systems, as can be seen in Figure 6d. All RhB@ZIF-71 samples in polar aprotic solvents exhibit a longer wavelength accompanied by a reduced intensity than in polar protic solvents. Remarkably, compared with the poor luminescence of RhB in highly polar aprotic solvents, the introduction of ZIF-71 greatly improves the luminescence of RhB (Figure S14). This phenomenon can be explained *via* the concept of aggregates we introduced above. When the pores of ZIF-71 contain monomers, there is enough room for the monomers to contact the aprotic solvent molecules to convert to lactone; when the pores contain aggregates (especially the J-type), it is likely that the remaining space in the pores cannot accommodate the solvent molecules. In other words, the ZIF-71 (host) can protect or shield the aggregates from direct exposure to the solvent molecules. On this basis, compared with the protic solvents, much more monomers inside the ZIF-71 pores convert to the colorless lactone in aprotic solvents, but the luminescence of aggregates is better protected. Consequently, RhB@ZIF-71 in polar aprotic solvents can exhibit luminescence with a longer wavelength and relatively lower intensity. This is the reason why RhB@ZIF-71 can show a better luminescence and sensing performance than pure RhB alone. Additionally, our interpretation can also be confirmed by observing the luminescent response of the RhB@ZIF-71 subject to polar protic solvents (Figure S14c), e.g., methanol (MeOH), ethanol (EtOH), and isopropanol (IPA) because Figure S14c reveals that the emission wavelength of RhB@ZIF-71 changes less than pure RhB, especially for IPA, which again can be attributed to the ZIF-71's protection of the RhB aggregates.

### 3. CONCLUSIONS

In summary, through a facile guest–host nanoconfinement strategy performed under ambient conditions, we demonstrate the encapsulation of fluorescent RhB monomers (or switchable aggregates) caged within the pores of ZIF-71. The new RhB@ZIF-71 system not only allows RhB to easily yield luminescence in the solid state but also provides remarkable mechanochromism, thermochromism, and solvatochromism properties that are not achievable to date by a traditional use of RhB dispersion in the liquid state or indeed by any other means of RhB@MOF systems known thus far. Above all, in the process of analyzing mechanochromism, thermochromism, and solvatochromism, we found several unique mechanisms



**Figure 5.** (a) Normalized excitation spectra (measured under  $\text{em}@650\text{ nm}$ ) and (b) normalized emission spectra (measured under  $\text{ex}@365\text{ nm}$ ) of RhB@ZIF-71 at different temperatures. (c) XRD patterns of RhB@ZIF-71 pellets after being tested at different temperatures. (d) Linear relationship of the emission peak wavelength as a function of temperature for RhB@ZIF-71.

summarized below. (i) Under mechanical stress, the RhB@ZIF-71 crystals deform by shear causing the conversion of H-type aggregates into J-type aggregates inside the MOF pores, giving rise to pressure sensing. (ii) The J-type aggregates are less affected by temperature due to the strong caging effect provided by the ZIF-71 pores, resulting in the red shift of emission in a very linear fashion, giving rise to noninvasive temperature sensing. (iii) The protective effect of ZIF-71 pores reduces the influence that solvents have on the RhB aggregates, leading to solvatochromic sensing of volatile organic compounds—previously undetected by unconfined RhB dyes alone.

Our RhB@ZIF-71 nanoconfinement strategy demonstrates the concept of a Guest@MOF system with multimodal fluorescent sensing response, presenting a new platform for the design of smart luminescent sensors. Moreover, we show that the exploitation of luminescent aggregates confined in MOF pores studied through the characterization of excitation spectrum combined with the luminescent lifetime data is a powerful approach for understanding the mechanisms of novel fluorochromic materials like the dye-encapsulated composite exemplified in this study.

## 4. EXPERIMENTAL SECTION

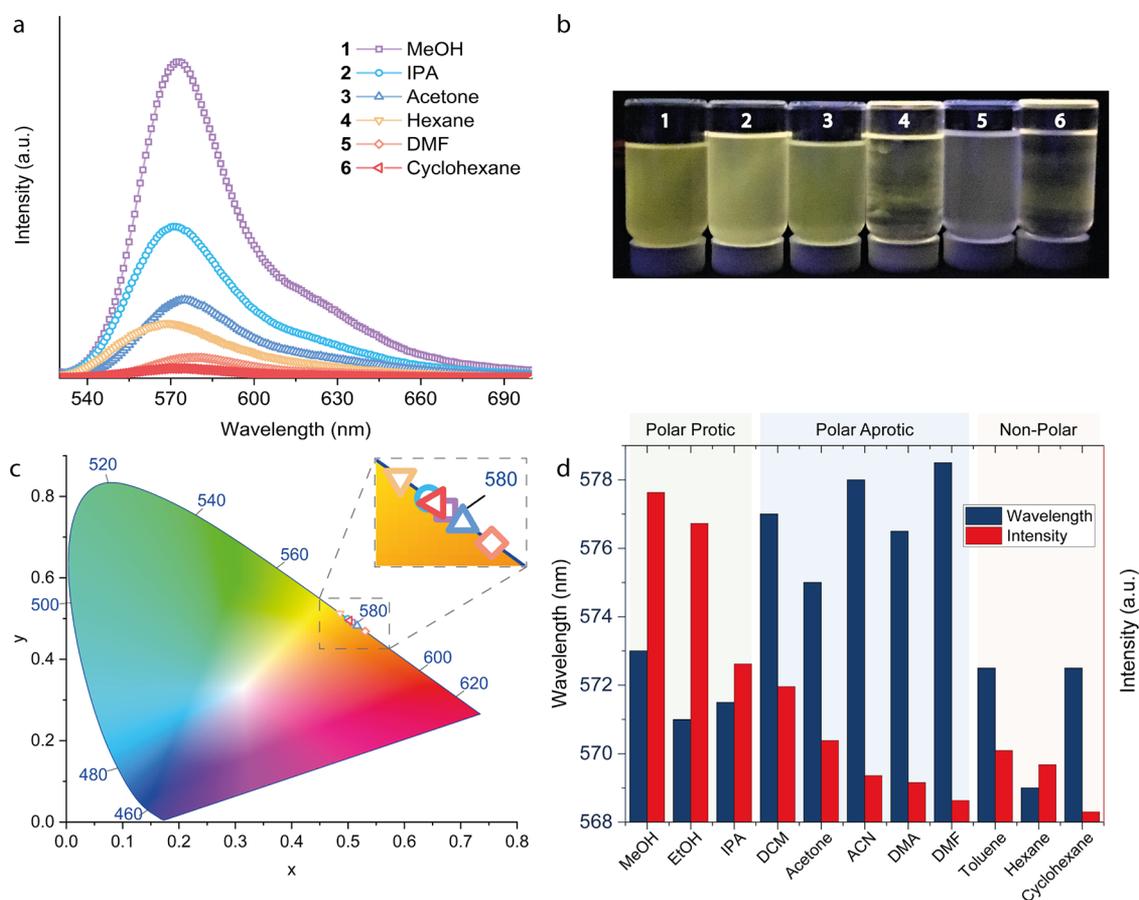
**4.1. Synthesis of RhB@ZIF-71.** A 90 mL methanol clear solution of 2.40 mmol of zinc acetate was rapidly poured into a 90 mL methanol solution of 9.60 mmol of dClm and different amounts of RhB (I, 0.01 mmol; II, 0.05 mmol; III, 0.5 mmol) under stirring. The mixed solution immediately changed from clear to turbid. After 24 h of stirring at room temperature, the sample was centrifuged at 8000 rpm to remove the excess reactants and subsequently washed twice with methanol (the sample was first put into methanol followed by sonication for 10 min, and then, the product was separated by

centrifugation) to remove the excess reactants and any RhB adhered on the surface of ZIF-71. We observed that, for the RhB<sub>I</sub>@ZIF-71 (0.01 mmol) and RhB<sub>II</sub>@ZIF-71 (0.05 mmol) samples, the methanol became clear after the second washing cycle, indicating the negligible loss of RhB guest molecules from the ZIF-71 host. Conversely, the RhB<sub>III</sub>@ZIF-71 (0.5 mmol) sample has RhB molecules that remained on the surface after two or more washing cycles, as evidenced from the pink color of the methanol solvent after multiple washes. The latter indicates surface adhesion of RhB molecules when a high concentration of the guest was being introduced during the synthesis step.

The procedure for preparing ZIF-71 was the same, except no RhB was added during the synthesis process.

**4.2. Sample Preparation for Fluorochromic Characterization.** For mechanochromism, the powders were pressed into pellets by using manual hydraulic pressure equipment with a 1 cm diameter die under forces of 1, 2, 3, and 4 tonnes. For thermochromism, no extra preparation was required. For solvatochromism, a concentration of 1 mg of RhB@ZIF-71 in 20 mL of solvents was chosen to avoid the negative influence of too intense emission. Then, to compare the performance of pure RhB and RhB@ZIF-71 in different solvents, several different concentrations of RhB-MeOH solutions were tested, and a concentration of  $7.5 \times 10^{-6}\text{ M}$  was selected at which the emission peak wavelength of pure RhB was identical to that of 1 mg of RhB@ZIF-71 in 20 mL of MeOH. Subsequently, the solutions of pure RhB in different solvents were also formulated at this concentration.

**4.3. Materials Characterization.** The structures and morphologies were examined under scanning electron microscopy (SEM; Carl Zeiss EVO LS15) and atomic force microscopy (AFM; Veeco Dimension 3100). An FS-5 spectrofluorometer (Edinburgh Instruments) was used to characterize the steady-state emission, excitation spectra, QY, CIE 1931, and lifetime measurements (more details are described in the Supporting Information). FTIR and Raman results were recorded by using a Nicolet iS10 FTIR spectrometer and a MultiRam FT-Raman spectrometer (Bruker), respectively. A UV-



**Figure 6.** (a) Emission spectra (ex@525 nm) of RhB@ZIF-71 when exposed to different volatile organic compounds (VOCs). (b) Solvatochromism observed under the 365 nm UV excitation where the concentration ratio used was 1 mg of RhB@ZIF-71 dispersed in 20 mL of the solvent. (c) Color variation presented on the CIE 1931 chromaticity diagram. (d) Change in the emission wavelength and peak intensity of RhB@ZIF-71 when tested in a wide range of polar protic, polar aprotic, and nonpolar solvents.

2600 UV–vis spectrophotometer (Shimadzu) was used to measure the absorption spectra and calculate the Kubelka–Munk (KM) function. The PXRD pattern was recorded using a Rigaku MiniFlex with a Cu  $K\alpha$  source (1.541 Å). TGA was performed using a TGA-Q50 machine (TA Instruments) equipped with a platinum sample holder under an  $N_2$  inert atmosphere at a heating rate of 10 °C/min from 50 to 800 °C.

## ■ ASSOCIATED CONTENT

### Supporting Information

The Supporting Information is available free of charge at <https://pubs.acs.org/doi/10.1021/acsami.0c10257>.

Analysis to estimate the chemical formula of RhB@ZIF-71, SEM images, X-ray diffraction patterns and crystallinity analysis, Raman spectra, excitation–emission spectra, lifetime emission spectra, spectrofluorometer characterization techniques, and a physically mixed sample of {RhB + ZIF-71} (PDF)

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## Notes

The authors declare no competing financial interest.

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